

**Eleven-dimensional symmetric supergravity
backgrounds, their geometric superalgebras,
and a common reduction**

Frank Klinker

Preprint 2014-06

October 2014

Eleven-dimensional symmetric supergravity backgrounds, their geometric superalgebras, and a common reduction

Frank Klinker^{1*}

Abstract. We present two different families of eleven-dimensional manifolds that admit non-restricted extensions of their isometry algebras to geometric superalgebras. Both families admit points for which the superalgebra extends to a super Lie algebra; on the one hand, a family of $N = 1, \nu = 3/4$ supergravity backgrounds and, on the other hand, an $N = 1, \nu = 1$ supergravity background. Furthermore, both families admit a point that can be identified with an $N = 4, \nu = 1/2$ six-dimensional supergravity background.

PACS codes: 04.65.+e, 02.40.Hw, 12.60.Jv

1 The setup

1.1 CW-spaces

In this text we consider CW-spaces and ask about conditions such that they can be considered as backgrounds in supergravity. In this context CW-spaces have been discussed in [1–6], for example. CW-spaces are Lorentzian solvable symmetric spaces that have been characterized in the early 1970's by M. Cahen and N. Wallach, see [7]. There is a one-to-one correspondence between $D = n + 2$ -dimensional CW-spaces and triples $(V, B, \langle \cdot, \cdot \rangle)$ of an n -dimensional euclidean vector space V , a symmetric map $B \in \text{End}(V)$, and an extension $\langle \cdot, \cdot \rangle$ of the euclidean product on V to a block-diagonal Lorentzian metric on $W = \mathbb{R}^2 \oplus V$. We fix a null basis $\{e_+, e_-\}$ of the \mathbb{R}^2 -factor. The manifold structure of the CW-space is defined by a Lie algebra structure on $V^* \oplus W$ of which the non-vanishing brackets are given by

$$[v^*, e_-] = Bv, [v^*, w] = -\langle Bv, w \rangle e_+, [e_-, w] = w^*. \quad (1)$$

Here, $*$: $V \rightarrow V^*$ is defined by $\langle \cdot, \cdot \rangle$. In particular, two such spaces $M_B, M_{\bar{B}}$ are isometric if and only if there exists an orthogonal map A and a positive scalar β such

* Faculty of Mathematics, TU Dortmund University, 44221 Dortmund, Germany
frank.klinker@math.tu-dortmund.de

that $\tilde{B} = \beta A B A^t$, see [7, 8]. The CW-space M_B that is defined by such triple is indecomposable if and only if B is non-degenerate. This can for example be seen by taking a look at the explicit form of the metric (2).

1.2 The metric and the Killing vector fields

Using the exponential map $\mu(x) = \exp(x^+ e_+) \exp(x^- e_-) \exp(\sum_i x^i e_i)$ we get

$$\begin{aligned} \mu^{-1} d\mu &= \sum_i x^- dx^- \otimes e_i^* \\ &+ dx^- \otimes e_- + \sum_i dx^i \otimes e_i + (dx^+ - \frac{1}{2} \sum_{ij} B_{ij} x^i x^j dx^-) \otimes e_+ \end{aligned}$$

such that the local metric of the $D = n + 2$ -dimensional CW-space M_B is

$$g = 2dx^+ dx^- - \sum_{ij} B_{ij} x^i x^j (dx^-)^2 + \sum_i (dx^i)^2. \quad (2)$$

We can and will consider an orthonormal basis $\{e_i\}$ of V such that the symmetric map B is diagonal, namely $B = \text{diag}(\lambda_1^2, \dots, \lambda_n^2)$. If $\lambda_1^2 = \dots = \lambda_{r_1}^2 < \lambda_{r_1+1}^2 = \dots = \lambda_{r_2}^2 < \dots < \lambda_{r_{\hat{n}-1}+1}^2 = \dots = \lambda_n^2$ we write $\{1, \dots, n\} = \bigcup_{\alpha=1}^{\hat{n}} I_\alpha$ with $I_\alpha = \{r_{\alpha-1}, \dots, r_\alpha\}$ and $r_0 = 1, r_{\hat{n}} = n$.

The Killing vector fields of this metric are given by

$$\begin{aligned} K_{(+)} &= -\partial_+, \quad K_{(-)} = -\partial_- \\ K_{(i)} &= \cos(\lambda_i x^-) \partial_i + \lambda_i \cos(x^-) x^i \partial_+, \quad i = 1, \dots, n, \\ K_{(i^*)} &= -\lambda_i \sin(\lambda_i x^-) \partial_i + \lambda_i^2 \cos(x^-) x^i \partial_+, \quad i^* = 1, \dots, n, \\ K_{(ij)} &= x^j \partial_i - x^i \partial_j, \quad (ij) \in I_1^2 \cup \dots \cup I_{\hat{n}}^2. \end{aligned} \quad (3)$$

We denote the Lie algebra that is spanned by the Killing vector fields by \mathcal{K}_0 , see also [1, 2].

1.3 Connections on the spinor bundle

We consider the irreducible spin bundle S over M_B . We denote the images of the basis $\{e_+, e_-, e_i\}$ under the spin representations by $\{\Gamma_+, \Gamma_-, \Gamma_i\}$. They obey the usual Clifford relation $\Gamma_A \Gamma_B + \Gamma_B \Gamma_A = -2g_{AB}$. The spin bundle splits into two subbundles $S = S_- \oplus S_+$ where the first and second summand are the -1 - and $+1$ eigenspaces of $\sigma := \frac{1}{2}[\Gamma_+, \Gamma_-]$, respectively. The two projections on the subbundles are given by $\sigma_\pm = -\frac{1}{2}\Gamma_\mp \Gamma_\pm$. Due to $\Gamma_+^2 = 0$ we have $\Gamma_+ : S_+ \rightarrow S_-$ and $S_+ = \ker \Gamma_+$. We denote the components of a section $\xi \in \mathcal{S}$ with respect to the above decomposition by $\xi = \xi_1 + \xi_2$.

$D = 11$ symmetric supergravity backgrounds

The Levi-Civita connection on M_B induces a connection on the spinor bundle. It is given by

$$\begin{aligned} \nabla_+ \xi &= \partial_+ \xi, & \nabla_i \xi &= \partial_i \xi \\ \nabla_- \xi &= \partial_- \xi + \frac{1}{2} \sum_i x^i \Gamma_{+B}(e_i) \xi. \end{aligned} \quad (4)$$

From [9] we know that any connection D on S that is compatible with the symmetric structure of M_B is described by a pair of elements (c, d) of the Clifford algebra of V . It is given by

$$\begin{aligned} D_+ \xi_1 &= \nabla_+ \xi_1, & D_+ \xi_2 &= \nabla_+ \xi_2, \\ D_- \xi_1 &= \nabla_- \xi_1 + c \xi_1, & D_- \xi_2 &= \nabla_- \xi_2 + d \xi_2, \\ D_i \xi_1 &= \nabla_i \xi_1 - \frac{1}{2} \Gamma_{+s_{c,d}}(e_i) \xi_2, & D_i \xi_2 &= \nabla_i \xi_2. \end{aligned} \quad (5)$$

with $s_{c,d}(v)$ given by

$$s_{c,d}(e_i) = c \Gamma_i - \Gamma_i d. \quad (6)$$

Remark 1. If the spinor bundle is reducible and of the form $S \otimes \mathbb{C}^N$ the result on the form of the connection remains true but with the parameters c, d taking their values in the tensor product of the Clifford algebra and $\mathfrak{gl}_N \mathbb{C}$.

The parallel spinors with respect to this connection are given by $\xi = (\xi_1, \xi_2)$ with

$$\begin{aligned} \xi_1 &= \exp(x^- c) \xi_1^0 + \frac{1}{2} \sum_i x^i \Gamma_{+s_{c,d}}(e_i) \exp(x^- d) \xi_2^0, \\ \xi_2 &= \exp(x^- d) \xi_2^0, \end{aligned} \quad (7)$$

where $\xi^0 = (\xi_1^0, \xi_2^0)$ is a constant spinor subject to the condition

$$(q_{c,d}(e_i) + B(e_i)) \xi_2 = 0 \quad (8)$$

with

$$q_{c,d}(e_i) = s_{c,d} \circ s_{c,d}(e_i) = c^2 \Gamma_i + \Gamma_i d^2 - 2c \Gamma_i d. \quad (9)$$

2 Geometric superalgebras

A geometric superalgebra of a CW-space M_B is an extension of \mathcal{K}_0 to a superalgebra $\mathcal{K}_0 \oplus \mathcal{K}_1$ with the following properties:

1. \mathcal{K}_1 is a subset of the space of the sections that are parallel with respect to a connection D on a spinor bundle S .
2. There exist a linear map $\mathcal{L} : \mathcal{K}_0 \rightarrow \text{End}(\mathcal{K}_1)$, such that $[\mathcal{L}_X, \mathcal{L}_Y] = \mathcal{L}_{[X, Y]}$.

3. There exist a bilinear symmetric map $\{\cdot, \cdot\} : \mathcal{K}_1 \times \mathcal{K}_1 \rightarrow \mathcal{K}_0$ such that $2\{\mathcal{L}_X \xi, \xi\} = [X, \{\xi, \xi\}]$.

A geometric superalgebra is called irreducible if S is as a Clifford module. In case that the spinor bundle is reducible with $S \otimes \mathbb{C}^N$ we call the superalgebra N -extended; sometimes we write $N = 1$ instead of irreducible. A geometric superalgebra is called non-restricted if the space \mathcal{K}_1 contains all parallel spinors, otherwise it is called restricted. We denote by ν the quotient of $\dim \mathcal{K}_1$ and $\dim S$. Then a superalgebra extension of a CW-space with flat connection is non-restricted if and only if $\nu = 1$. Nevertheless, a superalgebra extension with $\nu < 1$ can be restricted or a non-restricted.

A geometric superalgebra is called supersymmetry algebra if the extension actually is a super Lie algebra, i.e. if in addition

4. $\mathcal{L}_{\{\xi, \xi\}} \xi = 0$ is fulfilled for all $\xi \in \mathcal{K}_1$.

In this situation we also say that the underlying space admits geometric supersymmetry.

In our situation \mathcal{L} is the spinorial Lie derivative, see [10]. It is properly defined for Killing vector fields X and given by

$$\mathcal{L}_X \xi = \nabla_X \xi - \Gamma(\nabla X) \xi. \quad (10)$$

In particular, for this map the second part of 2. is satisfied.

From now on we restrict to $D = 11$ and $n = 9$.

2.1 A family of $N = 1, \nu = 1$ geometric superalgebras

We consider connection (5) with

$$c = \alpha \Gamma_{123}, d = \beta \Gamma_{123}. \quad (11)$$

These Clifford elements obey

$$q_{c,d}(e_i) = \begin{cases} (\alpha - \beta)^2 \Gamma_i & \text{for } i \in \{1, 2, 3\} \\ (\alpha + \beta)^2 \Gamma_i & \text{for } i \in \{4, \dots, 9\} \end{cases} \quad (12)$$

Therefore, (c, d) yields a flat connection on the CW-space M_B that is defined by

$$B = -\text{diag}((\alpha - \beta)^2 \mathbb{1}_3, (\alpha + \beta)^2 \mathbb{1}_6). \quad (13)$$

It is indecomposable if and only if $\alpha \neq \pm \beta$. We write $\lambda_j^2 = -(\alpha - \beta)^2$ for $j \in \{1, 2, 3\}$ and $\lambda_j^2 = -(\alpha + \beta)^2$ for $j \in \{4, \dots, 9\}$ and set $\lambda_j = -i(\alpha \pm \beta)$.

D = 11 symmetric supergravity backgrounds

If we consider

$$\mathcal{K}_1 = \left\{ \xi \in \mathcal{S} \left| \begin{array}{l} \xi = \xi(\xi_1^0, \xi_2^0) = \exp(\alpha x^- \Gamma_{123}) \xi_1^0 \\ \quad + \left(1 + \frac{1}{2} \sum_i \Gamma_+ x^i (\alpha \Gamma_{123} \Gamma_i - \beta \Gamma_i \Gamma_{123}) \right) \\ \quad \cdot \exp(\beta x^- \Gamma_{123}) \xi_2^0 \\ \xi_1^0, \xi_2^0 \text{ constant, } \sigma_- \xi_2^0 = \sigma_+ \xi_1^0 = 0 \end{array} \right. \right\} \quad (14)$$

In particular we have $\dim \mathcal{K}_1 = \dim S$ such that $\nu = 1$.

For $\xi = \xi(\xi_1^0, \xi_2^0) \in \mathcal{K}_1$ we have $\mathcal{L}_X \xi \in \mathcal{K}_1$; more precisely

$$\begin{aligned} \mathcal{L}_{K_{(+)}} \xi(\xi_1^0, \xi_2^0) &= \xi(0, 0), \\ \mathcal{L}_{K_{(-)}} \xi(\xi_1^0, \xi_2^0) &= \xi(-c\xi_1^0, -d\xi_2^0), \\ \mathcal{L}_{K_{(i)}} \xi(\xi_1^0, \xi_2^0) &= \xi\left(\frac{1}{2}\Gamma_+ s_{c,d} \xi_2^0, 0\right), \\ \mathcal{L}_{K_{(i^*)}} \xi(\xi_1^0, \xi_2^0) &= \xi\left(-\frac{1}{2}\Gamma_+ B(e_i) \xi_2^0, 0\right), \\ \mathcal{L}_{K_{(ij)}} \xi(\xi_1^0, \xi_2^0) &= \xi\left(\frac{1}{2}\Gamma_{ij} \xi_1^0, \frac{1}{2}\Gamma_{ij} \xi_2^0\right). \end{aligned} \quad (15)$$

To complete the geometric superalgebra we need the map $\{\cdot, \cdot\} : \mathcal{K}_1 \times \mathcal{K}_1 \rightarrow \mathcal{K}_0$. We will define $\{\xi, \xi\}$ for $\xi = \xi(\xi_1^0, \xi_2^0)$ by giving its projections onto the different directions of \mathcal{K}_0 . The full map is then given by polarization. We consider the charge conjugation C in eleven dimensions. It is skew-symmetric and obeys

$$C(\xi, \eta) = C(\xi_1, \eta_2) - C(\eta_1, \xi_2).$$

We write

$$\begin{aligned} \{\xi, \xi\} &= \{\xi, \xi\}^+ K_{(+)} + \{\xi, \xi\}^- K_{(-)} + \sum_{i=1}^9 \{\xi, \xi\}^i K_{(i)} \\ &\quad + \sum_{i=1}^9 \{\xi, \xi\}^{i^*} K_{(i^*)} + \frac{1}{2} \sum_{i,j=1}^9 \{\xi, \xi\}^{ij} K_{(ij)} \end{aligned} \quad (16)$$

with

$$\begin{aligned} \{\xi, \xi\}^+ &= C(\xi_1^0, \Gamma_- \xi_1^0), \quad \{\xi, \xi\}^- = C(\xi_2^0, \Gamma_+ \xi_2^0), \quad \{\xi, \xi\}^i = 2C(\xi_2^0, \Gamma_i \xi_2^0), \\ \{\xi, \xi\}^{i^*} &= \frac{2i}{\lambda_i} C(\xi_1^0, \Gamma_{123} \Gamma_i \xi_2^0), \\ \{\xi, \xi\}^{ij} &= \begin{cases} -i\lambda_1 C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) & \text{for } i, j \in \{1, 2, 3\} \\ i\lambda_2 C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) & \text{for } i, j \in \{4, \dots, 9\} \end{cases} \end{aligned} \quad (17)$$

In fact, a calculation similar to those in [11] proves the compatibility of \mathcal{L} and $\{\cdot, \cdot\}$ and yields the following result. It is an extension of the results obtained in [2] and a special case of a more general classification result on supergravity backgrounds of Cahen-Wallach type of which a publication is in preparation.

Proposition 2. *The CW-space given by $B = -\text{diag}((\alpha - \beta)^2 \mathbb{1}_3, (\alpha + \beta)^2 \mathbb{1}_6)$ along with the connection D described in (5) with $(c, d) = (\alpha \Gamma_{123}, \beta \Gamma_{123})$ defines a 1-parameter family of non-restricted irreducible geometric superalgebra if the odd-odd-bracket is defined as in (17).*

2.2 A family of non-restricted $N = 1, \nu = 3/4$ geometric superalgebras

We follow the way of construction from the preceding section. For this we consider now a connection that is defined by a pair of Clifford elements

$$\begin{aligned} c &= (\alpha \Gamma_{12} + \beta \Gamma_{45}) \Gamma_3 = (\alpha_+ X_{1245}^+ + \alpha_- X_{1245}^-) \Gamma_{123} \\ d &= (\alpha' \Gamma_{12} + \beta' \Gamma_{45}) \Gamma_3 = (\alpha'_+ X_{1245}^+ + \alpha'_- X_{1245}^-) \Gamma_{123} \end{aligned} \quad (18)$$

where we introduced the combinations $\alpha_{\pm} = \alpha \mp \beta$, $\alpha'_{\pm} = \alpha' \mp \beta'$ and the projection operators $X_{1245}^{\pm} = \frac{1}{2}(\mathbb{1} \pm \Gamma_{1245})$. The quadratic map $q_{c,d}$ associated to these elements is given by

$$q_{c,d}(e_i) = \begin{cases} (\alpha_- - \alpha'_+)^2 \Gamma_i X_{1245}^+ + (\alpha_+ - \alpha'_-)^2 \Gamma_i X_{1245}^- & \text{for } i \in \{1, 2\}, \\ (\alpha_+ - \alpha'_+)^2 \Gamma_i X_{1245}^+ + (\alpha_- - \alpha'_-)^2 \Gamma_i X_{1245}^- & \text{for } i \in \{3\}, \\ (\alpha_- + \alpha'_+)^2 \Gamma_i X_{1245}^+ + (\alpha_+ + \alpha'_-)^2 \Gamma_i X_{1245}^- & \text{for } i \in \{4, 5\}, \\ (\alpha_+ + \alpha'_+)^2 \Gamma_i X_{1245}^+ + (\alpha_- + \alpha'_-)^2 \Gamma_i X_{1234}^- & \text{for } i \in \{6, 7, 8, 9\}. \end{cases} \quad (19)$$

Therefore, we consider the CW-space M_B with

$$B = -\text{diag}((\alpha_- - \alpha'_+)^2 \mathbb{1}_2, (\alpha_+ - \alpha'_+)^2 \mathbb{1}_1, (\alpha_- + \alpha'_+)^2 \mathbb{1}_2, (\alpha_+ + \alpha'_+)^2 \mathbb{1}_4). \quad (20)$$

It is indecomposable if and only if $\alpha_{\pm} \neq \pm \alpha'_{\pm}$. As before, we write $\lambda_j^2 = -(\alpha_- - \alpha'_+)^2$ for $j \in \{1, 2\}$, $\lambda_3^2 = -(\alpha_+ - \alpha'_+)^2$, $\lambda_j^2 = -(\alpha_- + \alpha'_+)^2$ for $j \in \{4, 5\}$, $\lambda_j^2 = -(\alpha_+ + \alpha'_+)^2$ for $j \in \{6, \dots, 9\}$, and set $\lambda_j = -i(\alpha_{\pm} \pm \alpha'_{\pm})$. The connection defined by (c, d) above is non-flat such that $\nu = 1$ cannot be obtained. The parallel spinors in this situation define \mathcal{K}_1 and are given by

$$\mathcal{K}_1 = \left\{ \xi \in \mathcal{S} \left| \begin{array}{l} \xi = \xi(\xi_1^0, \xi_2^0) = \exp(x^- c) \xi_1^0 \\ \quad + \left(1 + \frac{1}{2} \sum_i \Gamma_i x^i s_{c,d}(e_i)\right) \exp(x^- d) \xi_2^0 \\ \xi_1^0, \xi_2^0 \text{ constant, } \sigma_- \xi_2^0 = \sigma_+ \xi_1^0 = 0, X_{1245}^- \xi_2^0 = 0 \end{array} \right. \right\} \quad (21)$$

such that $\nu = 3/4$. Again we have $\mathcal{L}_X \xi \in \mathcal{K}_1$ for all $\xi = \xi(\xi_1^0, \xi_2^0) \in \mathcal{K}_1$ with the same relations as before, namely (15).

We also complete the structure by introducing the map $\{\cdot, \cdot\} : \mathcal{K}_1 \times \mathcal{K}_1 \rightarrow \mathcal{K}_0$ as

$D = 11$ symmetric supergravity backgrounds

before. In this case it is given by

$$\begin{aligned} \{\xi, \xi\}^+ &= C(\xi_1^0, \Gamma_- \xi_1^0), \quad \{\xi, \xi\}^- = C(\xi_2^0, \Gamma_+ \xi_2^0), \quad \{\xi, \xi\}^i = 2C(\xi_2^0, \Gamma_i \xi_2^0), \\ \{\xi, \xi\}^{i*} &= \frac{2i}{\lambda_i} C(\xi_1^0, \Gamma_{125} \Gamma_i \xi_2^0), \\ \{\xi, \xi\}^{ij} &= \begin{cases} i\lambda_1 C(\xi_2^0, \Gamma_+ \Gamma_3 \xi_2^0) & \text{for } (ij) = (12) \\ i\lambda_3 C(\xi_2^0, \Gamma_+ \Gamma_{12345} \xi_2^0) & \text{for } (ij) = (45) \\ i\lambda_6 C(\xi_2^0, \Gamma_+ \Gamma_{123ij} \xi_2^0) & \text{for } i, j \in \{6, \dots, 9\} \end{cases} \end{aligned} \quad (22)$$

The involved calculations that prove the compatibility of \mathcal{L} and $\{\cdot, \cdot\}$ can be found in [11].

We summarize the above in the following statement.

Proposition 3. *The CW-space that is given by the symmetric map $B = -\text{diag}((\alpha_- - \alpha'_+)^2 \mathbb{1}_2, (\alpha_+ - \alpha'_+)^2, (\alpha_- - \alpha'_+)^2 \mathbb{1}_2, (\alpha_+ + \alpha'_+)^2 \mathbb{1}_4)$ together with the connection D described in (5) with $(c, d) = ((\alpha_+ X_{1245}^+ + \alpha_- X_{1245}^-) \Gamma_{123}, \alpha'_+ X_{1245}^+ \Gamma_{123})$ defines a 3-parameter family of non-restricted irreducible geometric superalgebras with $\nu = 3/4$ if the odd-odd-bracket is defined as in (22).*

Remark 4. • The first three bracket projections in (17) and (22) are the analog of the usual supersymmetry brackets as known from the super Poincaré-algebra in the flat situation. In more common notation it reads as $\{Q_\alpha, Q_\beta\}^\mu = \Gamma_{\alpha\beta}^\mu$. The two further projections are strongly related to the ingredients that enter into the definition of the superalgebra, namely the coefficients of the connection that defines the odd summand.

- In the N -extended situation the charge conjugation is replaced by the tensor product of a charge conjugation on the first factor and a bilinear form on the second factor in the construction of the odd-odd bracket.

Remark 5. The parameters of the families of geometric algebras in Propositions 2 and 3 can be reduced by one if we identify isometric Cahen Wallach spaces, so we are left with a 1-parameter family and a 2-parameter family, respectively.

2.3 Supersymmetry algebras

Propositions 2 and 3 tell us what the geometric superalgebras look like. The next question we will discuss is: when does such algebra yield a supersymmetry algebra? Or: what CW-space can be considered as supergravity background? The obstruction to this is the cubic spinorial condition 4, namely $\mathcal{L}_{\{\xi, \xi\}} \xi = 0$. Using (16) this is

$$\begin{aligned} \mathcal{L}_{\{\xi, \xi\}} \vec{\xi} &= \{\xi, \xi\}^+ \mathcal{L}_{K_{(+)}} \xi + \{\xi, \xi\}^- \mathcal{L}_{K_{(-)}} \xi + \sum_i \{\xi, \xi\}^i \mathcal{L}_{K_{(i)}} \xi \\ &\quad + \sum_i \{\xi, \xi\}^{i*} \mathcal{L}_{K_{(i^*)}} \xi + \frac{1}{2} \sum_{ij}^* \{\xi, \xi\}^{ij} \mathcal{L}_{K_{(ij)}} \xi. \end{aligned} \quad (23)$$

F. Klinker

If we split this into its two components we see that it yields one cubic equation for the constant spinor ξ_2^0 and one cubic equation for the two constant spinors ξ_1^0, ξ_2^0 . For example, in case of the geometric superalgebra from Proposition 2 the two equations for the second and first components are

$$0 = \beta C(\xi_2^0, \Gamma_+ \xi_2^0) \Gamma_{123} \xi_2 + \frac{1}{4}(\alpha - \beta) \sum_{ij \in \{1,2,3\}} C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) \Gamma_{ij} \xi_2^0 - \frac{1}{4}(\alpha + \beta) \sum_{ij \in \{4, \dots, 9\}} C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) \Gamma_{ij} \xi_2^0 \quad (24)$$

and

$$0 = \alpha(\xi_2, \Gamma_+ \xi_2^0) \Gamma_I \xi_1 + (\alpha - \beta) \sum_{i \in \{1,2,3\}} (C(\xi_1^0, \Gamma_i \xi_2^0) \Gamma_{123} \Gamma_i \xi_2^0 - C(\xi_1^0, \Gamma_{123} \Gamma_i \xi_2^0) \Gamma_i \xi_2^0) + (\alpha + \beta) \sum_{i \in \{4, \dots, 9\}} (C(\xi_1^0, \Gamma_i \xi_2^0) \Gamma_{123} \Gamma_i \xi_2^0 - C(\xi_1^0, \Gamma_{123} \Gamma_i \xi_2^0) \Gamma_i \xi_2^0) + \frac{1}{4}(\alpha - \beta) \sum_{ij \in \{1,2,3\}} C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) \Gamma_{ij} \xi_1^0 - \frac{1}{4}(\alpha + \beta) \sum_{ij \in \{4, \dots, 9\}} C(\xi_2^0, \Gamma_{123} \Gamma_{ij} \xi_2^0) \Gamma_{ij} \xi_1^0. \quad (25)$$

The only combinations of coefficients for which (25) and (24) can be identically solved is $\alpha = -3\beta$.

The analogue of (24)-(25) for the superalgebras from Proposition 3 can be found in [11] and the only choice of coefficients that solve the resulting equations is $\alpha_+ = -3\alpha'_+$.

We collect the results in the following proposition.

Proposition 6. 1. *The geometric superalgebra according to Proposition 2 yields non restricted $\nu = 1$ geometric supersymmetry if and only if*

$$B = -4\alpha^2 \text{diag}(4\mathbb{1}_3, \mathbb{1}_6)$$

and (c, d) is given by

$$c = -3\alpha\Gamma_{123}, \quad d = \alpha\Gamma_{123}.$$

2. *The geometric superalgebra according to Theorem 3 yields non restricted $\nu = 3/4$ geometric supersymmetry if and only if*

$$B = -\text{diag}((\alpha_- - \alpha'_+)^2 \mathbb{1}_2, 16\alpha_+^{\prime 2} \mathbb{1}_1, (\alpha_- + \alpha'_+)^2 \mathbb{1}_2, 4\alpha_+^{\prime 2} \mathbb{1}_4)$$

$D = 11$ symmetric supergravity backgrounds

and (c, d) is given by

$$c = (-3\alpha'_+ X_{1245}^+ + \alpha_- X_{1245}^-) \Gamma_{123}, \quad d = \alpha'_+ X_{1245}^+ \Gamma_{123}.$$

Remark 7. Part one of the above Proposition is exactly the unique maximal supergravity background of CW-type that has been considered in [1, 2]. The uniqueness follows after identifying isometric spaces, see Remark 5.

3 A common reduction

Consider two CW-spaces of dimension $D' = n' + 2$ and $D = n + 2 > D'$ associated to the symmetric maps B' and $B = B' \oplus B''$. These two spaces come with geometric superalgebras \mathcal{K}' and \mathcal{K} that are N' -extended and N -extended. Furthermore the relation $N' = 2 \lfloor \frac{n-n'}{2} \rfloor N$ holds.

Then \mathcal{K}' is called a reduction of \mathcal{K} , or \mathcal{K} an oxidation of \mathcal{K}' , if the following holds: We erase from \mathcal{K}_0 exactly $n - n'$ Killing vector fields from $K_{(i)}$ and the same from $K_{(i^*)}$ such that the remaining part of \mathcal{K} remains an algebra and is isomorphic to \mathcal{K}' – maybe after restricting \mathcal{K}_1 . This is in particular of interest in the following situations:

- Can we reduce and/or oxidate a geometric supersymmetry so that the result is such, too?
- Suppose no member of a family of geometric superalgebras yields geometric supersymmetry, can we find a reduction of a member that does?

We will address these question with regard to the families we presented before.

3.1 The first reduction

Due to the nature of the connection defined by (18) the odd part of the super Lie algebra from Proposition 6-2 cannot be restricted further in the generic situation. Nevertheless, there is one configuration of parameters where this is possible, namely $\alpha'_+ = 0$. In this situation, we are left with a decomposable CW-space associated to $B = -\alpha^2 \text{diag}(1, 1, 0, 1, 1, 0, 0, 0, 0)$ and the odd part of the geometric supersymmetry is restricted to $X_{1245}^+ \xi_1^0 = 0$ in addition to $X_{1245}^- \xi_2^0 = 0$, i.e. $\nu = \frac{1}{2}$.

By taking a closer look at (24)-(25) we see that both terms vanish for a truncated summation over $\{1, 2, 4, 5\}$ if the proposed restriction is performed. The resulting algebra can then be interpreted as a reduced 4-extended geometric supersymmetry in the following way. These are exactly the data for the $D = 6$, $N = 4$ supergravity background proposed in [12].

Consider the six-dimensional CW-space associated to $B = -\alpha^2 \mathbb{1}_4$ and its spinor bundle S_6 with charge conjugation C_6 . Within $S = S_6 \otimes \mathbb{C}^4$ we identify the second factor with the five-dimensional spin-representation and provide it with the charge

conjugation C_5 . Then $C = C_6 \otimes C_5$ defines a bilinear form on S . Furthermore we consider the connection defined by

$$c = \alpha X_{1234}^- \Gamma_{12} \otimes T, \quad d = 0. \quad (26)$$

Here T is some vector in \mathbb{R}^5 regarded as an element of the Clifford algebra with $T^2 = -\mathbb{1}$. The parallel spinors of this connection are parametrized by constant spinors $\xi_1^0 \otimes v_1 \in S_{6,-} \otimes \mathbb{C}^5$, $\xi_2^0 \otimes v_2 \in X_{1234}^+ S_{6,+} \otimes \mathbb{C}^5$.

To define \mathcal{K}_1 we impose the further condition $X_{1234}^+ \xi_1^0 \otimes v_1 = 0$ that yields the reduction to $\nu = 1/2$.

We use C to write down the bracket structure of the algebra:

$$\begin{aligned} \{\xi, \eta\}^+ &= C_6(\xi_1^0, \Gamma_- \eta_1^0) C_5(v_1, w_1), \\ \{\xi, \eta\}^- &= C_6(\xi_2^0, \Gamma_+ \eta_2^0) C_5(v_2, w_2), \\ \{\xi, \eta\}^i &= C_6(\xi_1^0, \Gamma_i \eta_2^0) C_5(v_1, w_2) + C_6(\eta_1^0, \Gamma_i \xi_2^0) C_5(w_1, v_2), \\ \{\xi, \eta\}^{i*} &= \frac{1}{\alpha} (C_6(\xi_1^0, \Gamma_{12} \Gamma_i \eta_2^0) C_5(v_1, Tw_2) + C_6(\eta_1^0, \Gamma_{12} \Gamma_i \xi_2^0) C_5(w_1, Tw_2)), \\ \{\xi, \eta\}^{ij} &= -\alpha C_6(\xi_2^0, \Gamma_+ \eta_2^0) C_5(v_2, Tw_2) \end{aligned} \quad (27)$$

for $\xi = \xi(\xi_1^0 \otimes v_1, \xi_2^0 \otimes v_2)$, $\eta = \eta(\eta_1^0 \otimes w_1, \eta_2^0 \otimes w_2)$.

Proposition 8. *The above data yield a six dimensional 4-extended non-restricted geometric supersymmetry with $\nu = 1/2$ that is the reduction of the eleven-dimensional geometric supersymmetry from Proposition 6-2 for $\alpha'_+ = 0$.*

3.2 The second reduction

As we know, the non-restricted geometric superalgebras with $\nu = 1$ from Proposition 2 only yield geometric supersymmetry for a special choice of coefficients, see Proposition 6-1. Nevertheless, if we again take a look at (24)-(25) for the choice $\beta = 0$ we see that the truncation to $\{1, 2, 4, 5\}$ annihilates both sums if we consider a restriction to $\nu = 1/2$ of the odd part that is similar to the one before, namely $X_{1245}^+ \xi_1^0 = X_{1245}^- \xi_2^0 = 0$. The interpretation of this truncation as a reduction is as follows:

We consider the same six-dimensional CW-space and the same spinor bundle as before but with connection defined by

$$c = \alpha \Gamma_{12} \otimes T, \quad d = 0. \quad (28)$$

This connection is flat and the parallel spinors are parametrized by all of $\xi_1^0 \otimes v_1 \in S_{6,-} \otimes \mathbb{C}^4$, $\xi_2^0 \otimes v_2 \in S_{6,+} \otimes \mathbb{C}^4$. We consider the restriction subject to the conditions $X_{1234}^+ \xi_1^0 = X_{1234}^- \xi_2^0 = 0$ as before. Furthermore, we use the same brackets as in (27).

$D = 11$ symmetric supergravity backgrounds

Proposition 9. *The six-dimensional data above yield a restricted 4-extended geometric supersymmetry with $\nu = 1/2$ that is a reduction of the eleven-dimensional non-restricted geometric superalgebra from Proposition 2 for $\beta = 0$.*

3.3 Concluding remarks

- The correspondences that have been claimed in Propositions 8 and 9 can be made precise by identifying T with Γ_3 and embedding M_6 into M_{11} by $(\pm, 1, 2, 3, 4) \rightarrow (\pm, 1, 2, 4, 5)$.
- We will briefly explain why the two six-dimensional supersymmetries constructed in the last two subsections are essentially the same although the connections that define the structures are not.
In the construction of both supersymmetries we could have forgotten about the further restriction of ξ_1^0 . That would also lead to supersymmetries and to six-dimensional supergravity backgrounds but with $\nu = 3/4$ in both cases. The first one would then be non-restricted and the second one would be restricted, by definition.
However, by introducing the further condition on ξ_1^0 we guarantee that the algebra structures in both cases coincide. Roughly, this is due to the fact that both connections differ by a half spinor projection that enters into the bracket structure (27). In case of additional restriction this is only artificially present so that we could omit it.
- We want to emphasize the differences of the two oxidations of Propositions 8 and 9.
In the first case the eleven-dimensional CW space is decomposable with non-flat connection and the geometric superalgebra is indeed a super Lie algebra. Therefore, oxidation and reduction can both be considered as supergravity backgrounds.
In the second case the oxidation is an indecomposable CW-space with flat connection but the superalgebra does not define supersymmetry. Nevertheless, the space belongs to a family of geometric superalgebras that contains a supersymmetric solution, namely the sole maximal supergravity background of CW-type, see Remark 7.

References

- [1] Piotr T. Chruściel and Jerzy Kowalski-Glikman (1984) *Phys. Lett. B* **149** 107-110.
- [2] José Figueroa-O'Farrill and George Papadopoulos (2001) *J. High Energy Phys.* 2001 **08** 036.
- [3] Gerome P. Gauntlett and Christopher M. Hull (2002) *J. High Energy Phys.* 2002 **06** 013.
- [4] José Figueroa-O'Farrill (2003) *Class. Quant. Grav.* **20** 3327-3340.
- [5] José Figueroa-O'Farrill, Patrick Meessen, and Simon Philip (2005) *Class. Quant. Grav.* **22** 207-226.

F. Klinker

- [6] José Figueroa-O'Farrill and Noel Hustler (2012) *Class. Quant. Grav.* **30** 045008.
- [7] Michel Cahen and Nolan Wallach (1970) *Bull. Amer. Math. Soc.* **76** 585-591.
- [8] Hung-Hsi Wu (1967) *Pacific J. Math.* **20** no. 2 351-392.
- [9] Frank Klinker (2014) *Adv. Appl. Clifford Algebr.* **24** no. 3 737-768.
- [10] Yvette Kosmann (1971) *Ann. Mat. Pura Appl.* **91** no. 1 317-395.
- [11] Frank Klinker (2014) *submitted*, Preprint: [arXiv:1406.4672](https://arxiv.org/abs/1406.4672) [math.DG].
- [12] Patrick Meessen (2002) *Phys. Rev. D* **65** 087501.

Preprints ab 2012/06

- 2014-06 **Frank Klinker**
Eleven-dimensional symmetric supergravity backgrounds, their geometric superalgebras, and a common reduction
- 2014-05 **Tomáš Dohnal and Hannes Uecker**
Bifurcation of nonlinear Bloch waves from the spectrum in the Gross-Pitaevskii equation
- 2014-04 **Frank Klinker**
A family of non-restricted $D = 11$ geometric supersymmetries
- 2014-03 **Martin Heida and Ben Schweizer**
Non-periodic homogenization of infinitesimal strain plasticity equations
- 2014-02 **Ben Schweizer**
The low frequency spectrum of small Helmholtz resonators
- 2014-01 **Tomáš Dohnal, Agnes Lamacz, Ben Schweizer**
Dispersive homogenized models and coefficient formulas for waves in general periodic media
- 2013-16 **Karl Friedrich Siburg**
Almost opposite regression dependence in bivariate distributions
- 2013-15 **Christian Palmes and Jeannette H. C. Woerner**
The Gumbel test and jumps in the volatility process
- 2013-14 **Karl Friedrich Siburg, Katharina Stehling, Pavel A. Stoimenov, Jeannette H. C. Wörner**
An order for asymmetry in copulas, and implications for risk management
- 2013-13 **Michael Voit**
Product formulas for a two-parameter family of Heckman-Opdam hypergeometric functions of type BC
- 2013-12 **Ben Schweizer and Marco Veneroni**
Homogenization of plasticity equations with two-scale convergence methods
- 2013-11 **Sven Glaser**
A law of large numbers for the power variation of fractional Lévy processes
- 2013-10 **Christian Palmes and Jeannette H. C. Woerner**
The Gumbel test for jumps in stochastic volatility models
- 2013-09 **Agnes Lamacz, Stefan Neukamm and Felix Otto**
Moment bounds for the corrector in stochastic homogenization of a percolation model
- 2013-08 **Frank Klinker**
Connections on Cahen-Wallach spaces
- 2013-07 **Andreas Rätz and Matthias Röger**
Symmetry breaking in a bulk-surface reaction-diffusion model for signaling networks
- 2013-06 **Gilles Francfort and Ben Schweizer**
A doubly non-linear system in small-strain visco-plasticity
- 2013-05 **Tomáš Dohnal**
Traveling solitary waves in the periodic nonlinear Schrödinger equation with finite band potentials

- 2013-04 **Karl Friedrich Siburg, Pavel Stoimenov and Gregor N. F. Weiß**
Forecasting portfolio-value-at-risk with nonparametric lower tail dependence estimates
- 2013-03 **Martin Heida**
On thermodynamics of fluid interfaces
- 2013-02 **Martin Heida**
Existence of solutions for two types of generalized versions of the Cahn-Hilliard equation
- 2013-01 **Tomáš Dohnal, Agnes Lamacz, Ben Schweizer**
Dispersive effective equations for waves in heterogeneous media on large time scales
- 2012-19 **Martin Heida**
On gradient flows of nonconvex functional in Hilbert spaces with Riemannian metric and application to Cahn-Hilliard equations
- 2012-18 **Robert V. Kohn, Jianfeng Lu, Ben Schweizer and Michael I. Weinstein**
A variational perspective on cloaking by anomalous localized resonance
- 2012-17 **Margit Rösler and Michael Voit**
Olshanski spherical functions for infinite dimensional motion groups of fixed rank
- 2012-16 **Selim Esedoğlu, Andreas Rätz, Matthias Röger**
Colliding Interfaces in Old and New Diffuse-interface Approximations of Willmore-flow
- 2012-15 **Patrick Henning, Mario Ohlberger and Ben Schweizer**
An adaptive multiscale finite element method
- 2012-14 **Andreas Knauf, Frank Schulz, Karl Friedrich Siburg**
Positive topological entropy for multi-bump magnetic fields
- 2012-13 **Margit Rösler, Tom Koornwinder and Michael Voit**
Limit transition between hypergeometric functions of type BC and Type A
- 2012-12 **Alexander Schnurr**
Generalization of the Blumenthal-Gettoor index to the class of homogeneous diffusions with jumps and some applications
- 2012-11 **Wilfried Hazod**
Remarks on pseudo stable laws on contractible groups
- 2012-10 **Waldemar Grundmann**
Limit theorem for radial random walks on Euclidean spaces of high dimensions
- 2012-09 **Martin Heida**
A two-scale model of two-phase flow in porous media ranging from porespace to the macro scale
- 2012-08 **Martin Heida**
On the derivation of thermodynamically consistent boundary conditions for the Cahn-Hilliard-Navier-Stokes system
- 2012-07 **Michael Voit**
Uniform oscillatory behavior of spherical functions GL_n/U_n at the identity and a central limit theorem
- 2012-06 **Agnes Lamacz and Ben Schweizer**
Effective Maxwell equations in a geometry with flat rings of arbitrary shape