

Experimental verification and comparison of analytical and FE models for calculation of a Bitter solenoid

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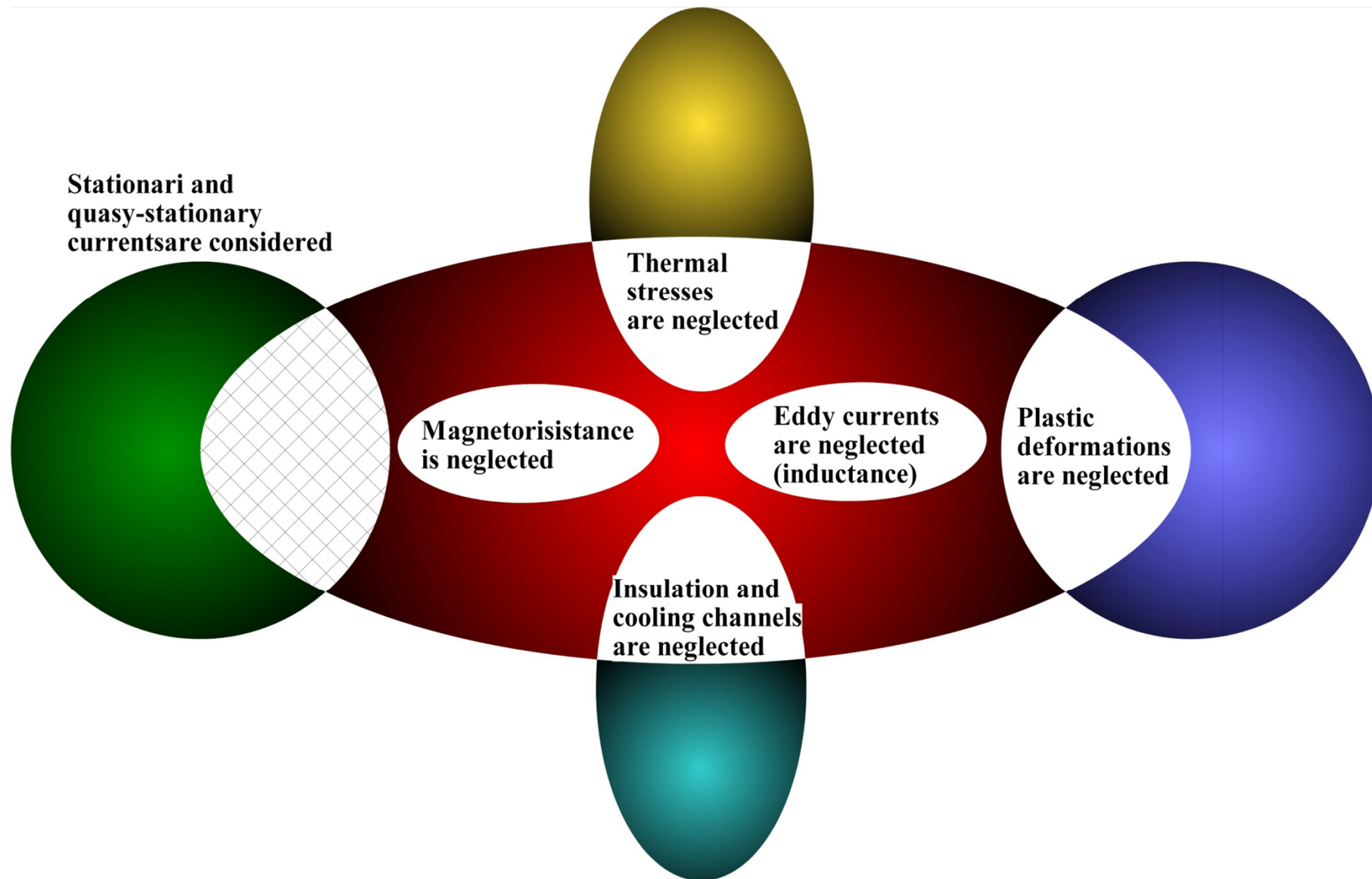
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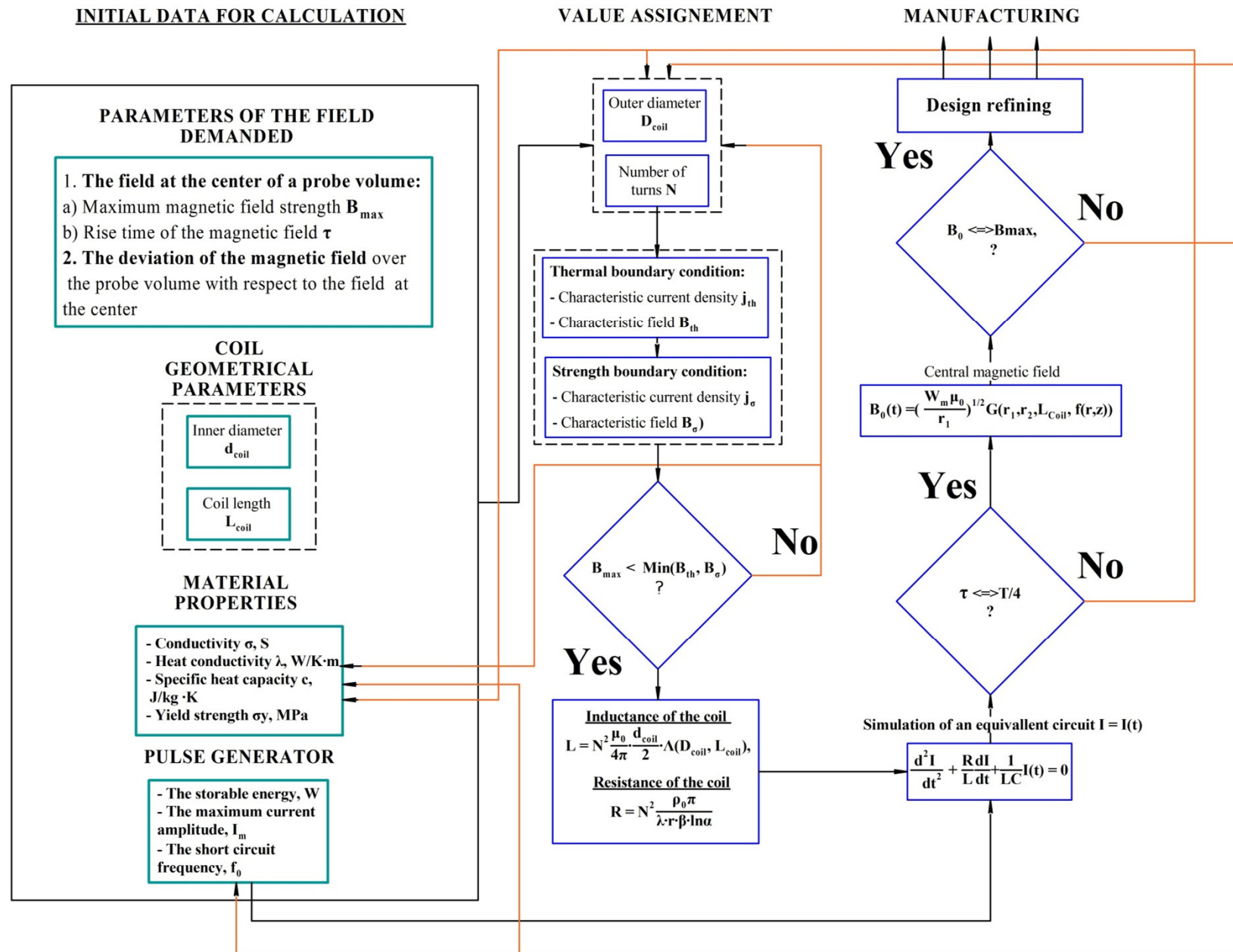
Analytical calculation of a Bitter solenoid

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1. Limitations of the approach



2. Calculation methodology



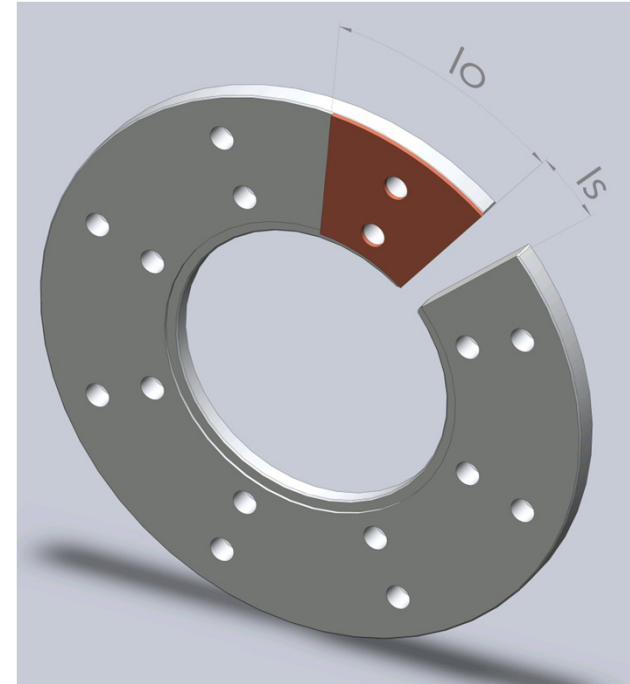
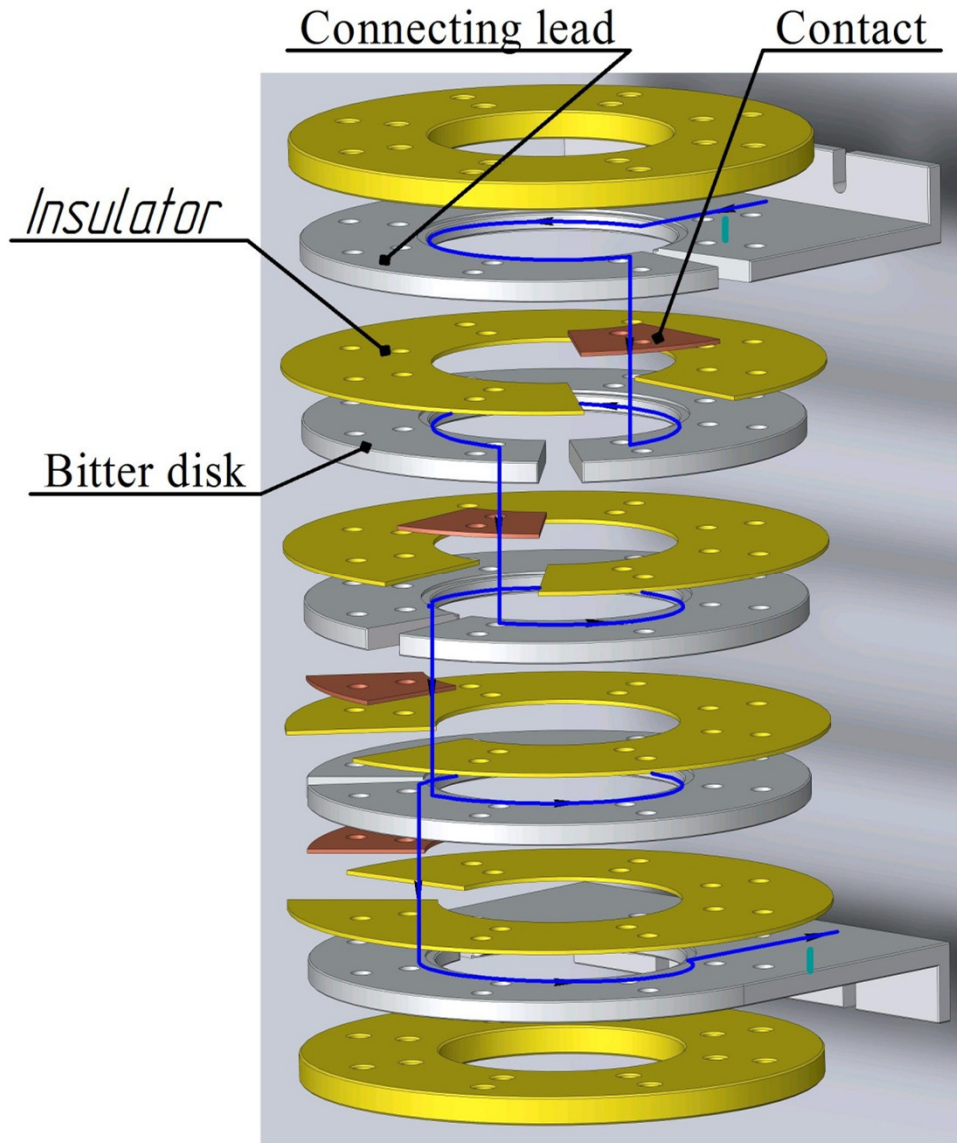
3. Determination of input data

1. Material properties:
 - a) Conductivity σ [MS]
 - b) Specific heat capacity c [J/kg·K]
 - c) Working temperatures: initial T_i and final T_f [K]
 - d) Mass density ρ [kg/m³]
 - e) Yield strength σ_y [MPa]

2. Value assignment for an outer radius r_2 and a number of turns N

3. To take into account a symmetry breakdown in a real coil

3. Determination of input data



Ideal symmetrical disk = 360°
Real disk = $360^\circ - l_o - l_s$



Nominal number of turns must be reduced to an effective ideal number of turns having a cylindrical symmetry.

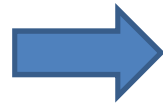
4. Boundary conditions: characteristic current densities and fields

1. Thermal BC takes into account only heating of the solenoid

Material integral must be as high as possible:

$$F_{mat}(T_i, T_f) := \int_{T_i}^{T_f} \rho_m \cdot \frac{c(T)}{\rho(T)} dT$$

Material	$F_{Mat}(T_i, T_f) [10^{16} A^2s/m^4]$	
	77 K – 400 K	77 K – 700 K
Cu	9,42	
Al	4,58	
C17510	4,45	
C17200	1,30	
AerMet 100	0,27	0,46
AISI 316	0,15	0,26



$$j_{th} := \sqrt{\frac{F_{mat}(T_i, T_f)}{t_{pulse} \cdot \zeta}}$$

$$B_{th} := \mu_0 \cdot \xi \cdot r_1 \cdot \sqrt{\frac{F(T_i, T_f)}{t_{pulse} \cdot \zeta}}$$

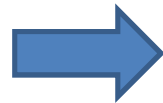
$$B_0 := \min(B_{th}, B_{\sigma}) \cdot \ln(\alpha)$$

2. Strength BC takes into account only mechanical strength of a material

$$\sigma_{max} = \sigma_y$$

ξ – filling factor of the coil

$\alpha = r_2/r_1$ – form – factor of the coil



$$j_{\sigma} := \sqrt{\frac{\sigma_{max}}{\mu_0}} \cdot \frac{1}{\xi \cdot r_1 \cdot \sqrt{\ln(\alpha)}}$$

$$B_{\sigma} := \sqrt{\mu_0 \cdot \sigma_{max}} \cdot \frac{1}{\sqrt{\ln(\alpha)}}$$

5. Calculation of the inductance and active resistance of the coil

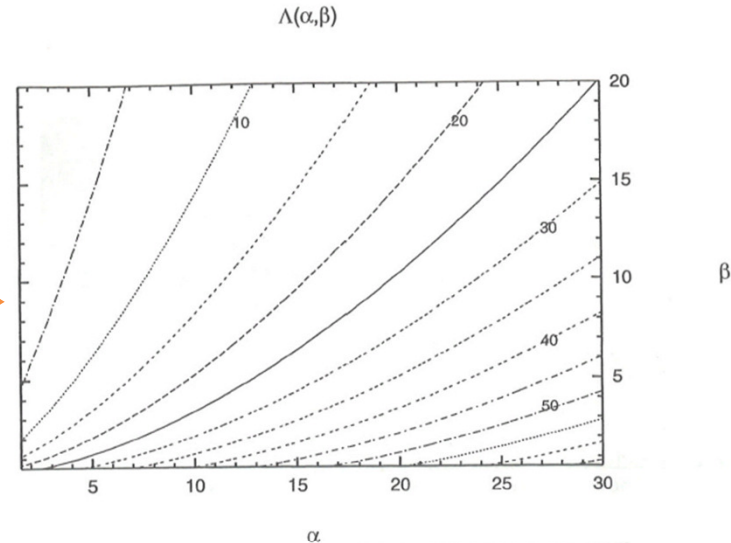
A) Inductance

$$L_{\text{coil}} := \eta \frac{2 \mu_0 \cdot r_1 \cdot \Lambda(\alpha, \beta)}{4 \cdot \pi}$$

$\beta = l_{\text{coil}}/2r_1$ – form – factor of the coil

η – effective number of turns

$\Lambda(\alpha; \beta)$ – self – inductance factor



B) Resistance

$$R_{\text{dc}} := \eta^2 \cdot \frac{\rho \cdot \pi}{\xi \cdot r_1 \cdot \beta \cdot \ln(\alpha)}$$

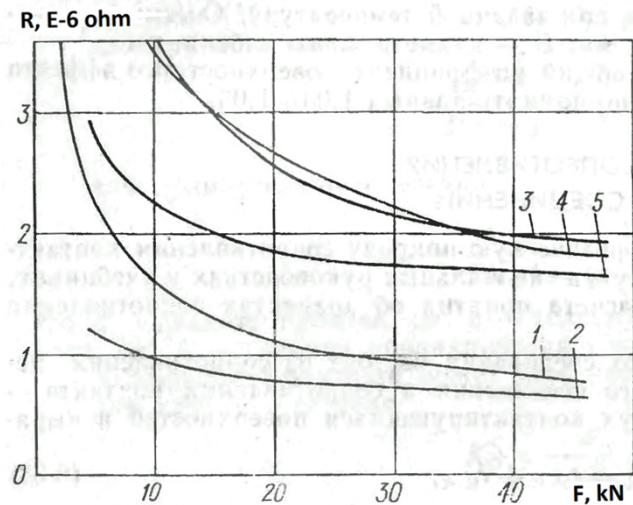
R_{contact}

$$R_{\text{ac}} := \frac{a \cdot (R_{\text{dc}} + R_{\text{contact}}) \cdot \sqrt{\omega d \cdot \frac{\mu_0}{2 \cdot \rho}}}{u}$$

a – cross-section of a turn [m²]

u – circumference of a turn [m]

ωd – current angular frequency [Hz]



- 1 – Cu-Cu contact, oxidized surface
- 2 – Cu-Cu contact, clean surface
- 3 – Contact between Cu-Cr-Zn plates, oiled surfaces,
- 4 – Cu-Cu, clean surface,
- 5 – Contact between Cu-Cr-Zn plates, clean surface,

6. Simulation of an equivalent RLC circuit

Differential equation of electromagnetic damped oscillation can be easily solved at initial conditions defined by a pulsed generator:

C – capacitance [F]

U – discharge voltage[V]

Given

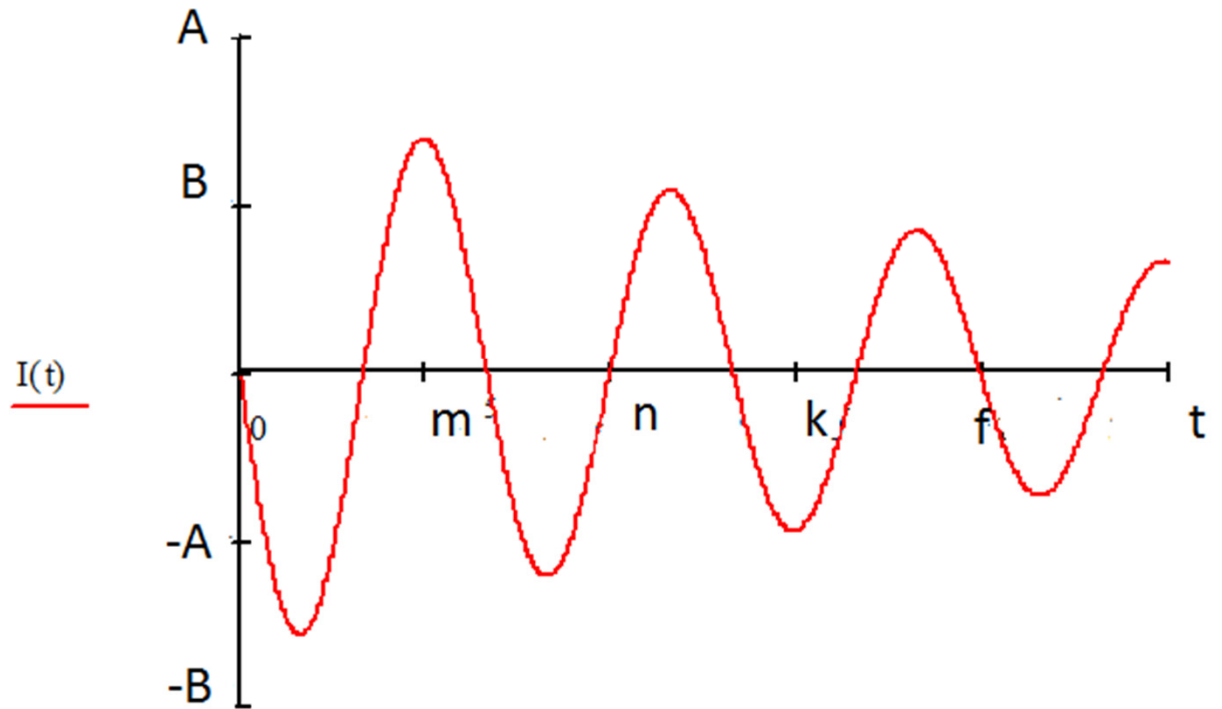
$$\frac{d^2}{dt^2}q(t) + 2 \cdot \beta_1 \cdot \left(\frac{d}{dt}q(t) \right) + \omega^2 \cdot q(t) = 0$$

$$q(0) = C \cdot U$$

$$q'(0) = 0$$

q := Odesolve(t, t1)

$$I(t) := \frac{d}{dt}q(t)$$



On this step a verification of the frequencies demanded and calculated takes place

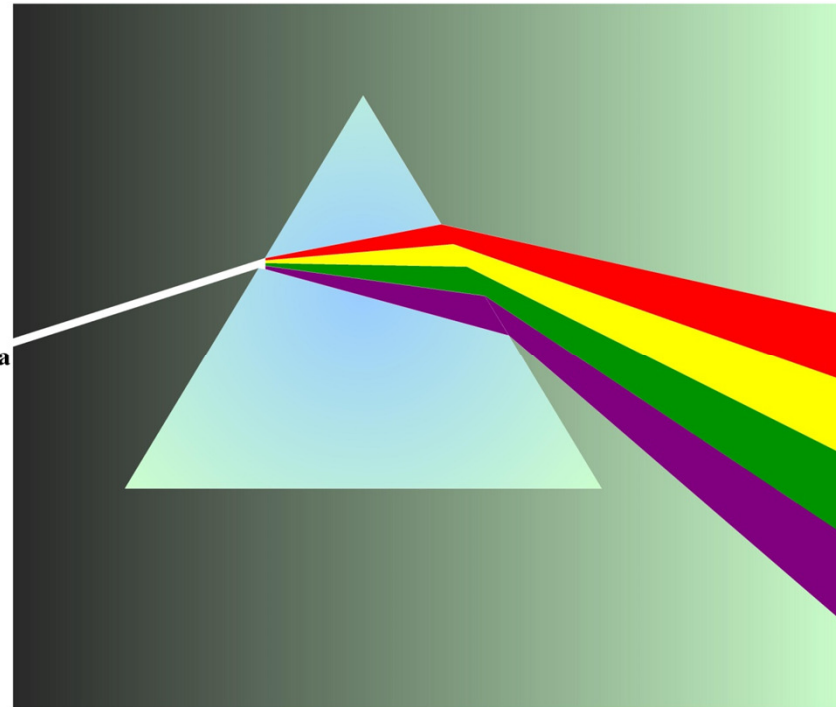
7. Calculation of the central magnetic field

$$f(r,z) := \frac{r1}{r} \quad \text{Current density distribution}$$

A – cross-section

$$B_0(t) := \frac{\sqrt{W_m(t) \cdot \frac{\mu_0}{r1}} \cdot \sqrt{2 \frac{\pi}{\Lambda(\alpha, \beta)}} \cdot \int_A \int \frac{r^2 \cdot f(r,z)}{(r^2 + z^2)^{\frac{3}{2}}} dr dz}{\int_A \int f(r,z) dr dz}$$

Fabry Formula



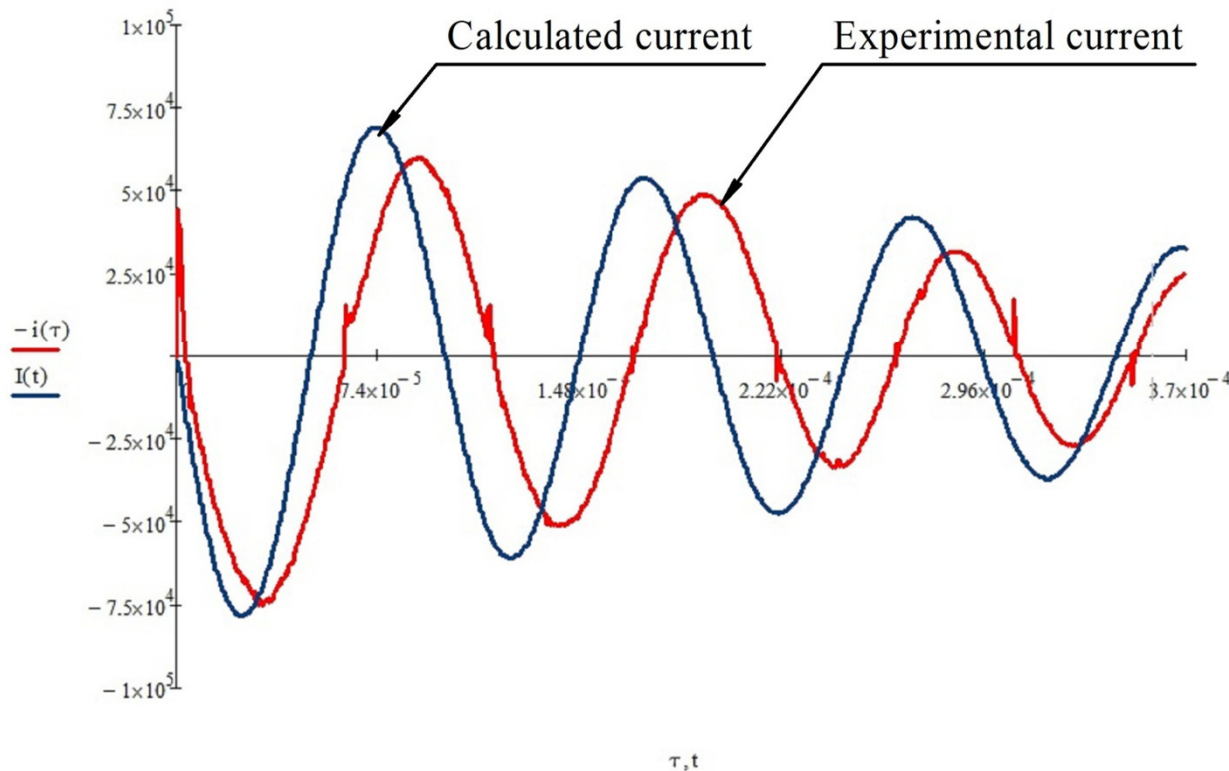
- The central Field B
- The magnetic energy Wm
- The Fabry factor G(α,β)
- The inner radius r

After the integration a closed form expression was obtained so a convenient calculation can be made for any sizes and current courses:

$$B_0(t) := \frac{\sqrt{W_m(t) \cdot \frac{\mu_0}{r1}} \cdot \ln \left(\frac{l_{coil}}{r1} + \sqrt{\frac{l_{coil}^2}{r1^2} + 1} \right)}{l_{coil} \cdot \ln \left(\frac{r2}{r1} \right)}$$

8. Verification of the methodology on an example of BWI Bitter solenoid

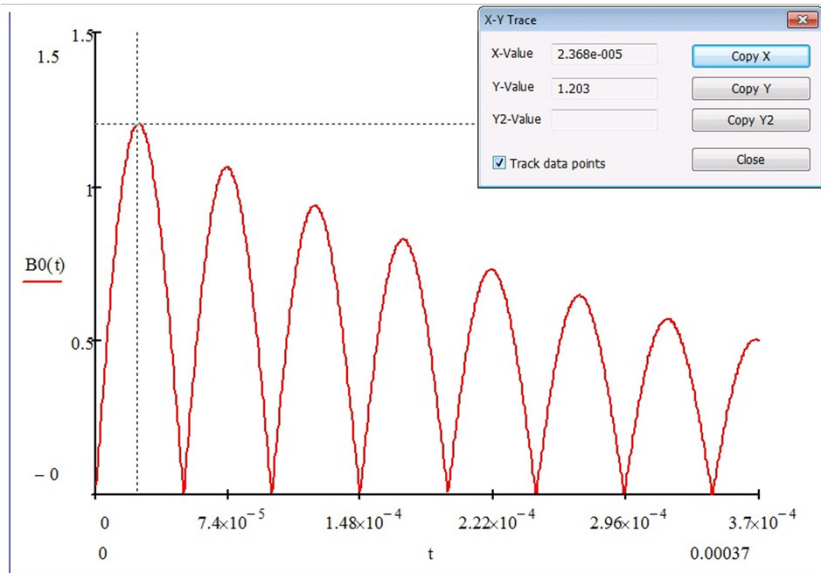
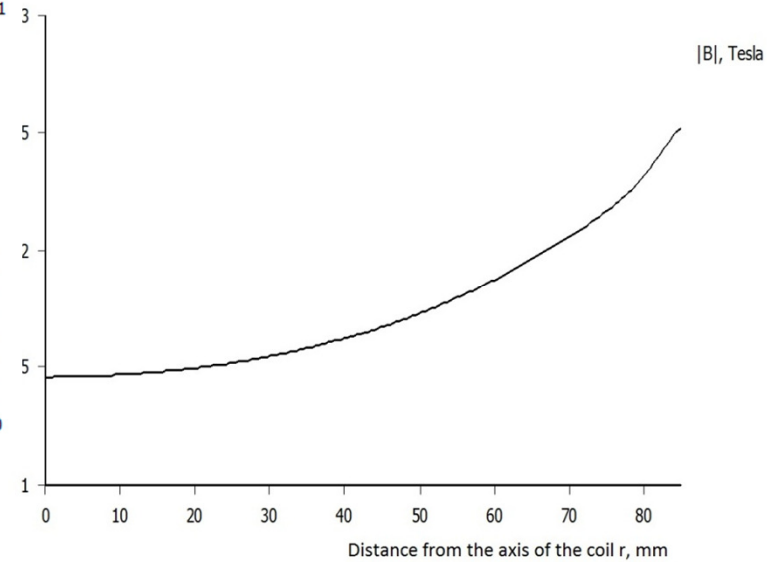
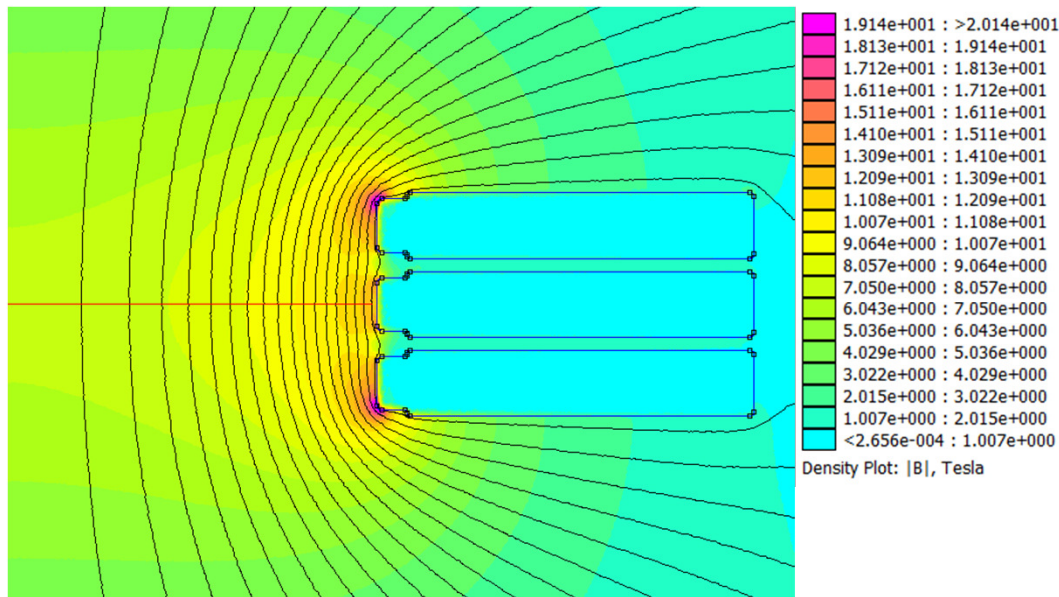
The initial conditions $C = 130e-6$ F, $U = 10$ kV and $I(0)=0$ correspond to the experiment:



Parameter	$R \cdot 10e-3$ [Ohm]	$L \cdot 10e-9$ [H]
Experimental	7.441	1504
Calculated	6.588	1838
Relative error	11.5%	22%

A comparison of the graphs allows to draw a conclusion of satisfying errors between the experimental and calculated curves.

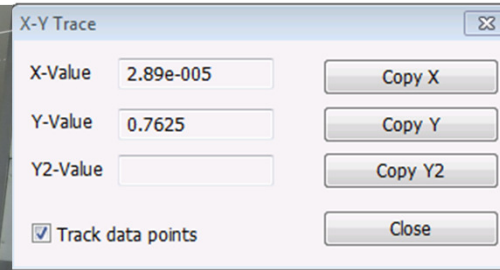
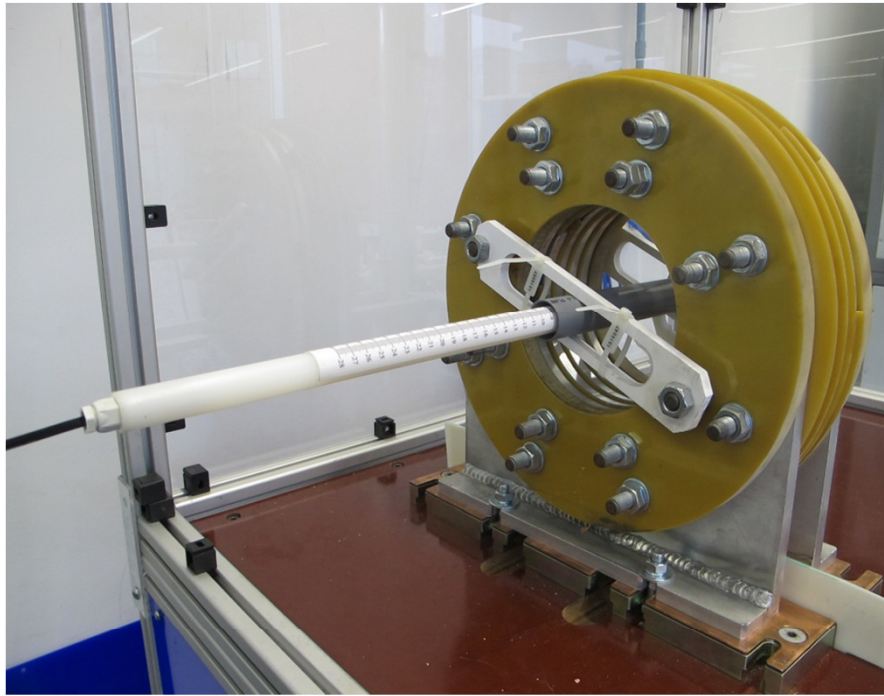
9. Modeling of the magnetic field in the probe volume of the BWI coil using FEMM



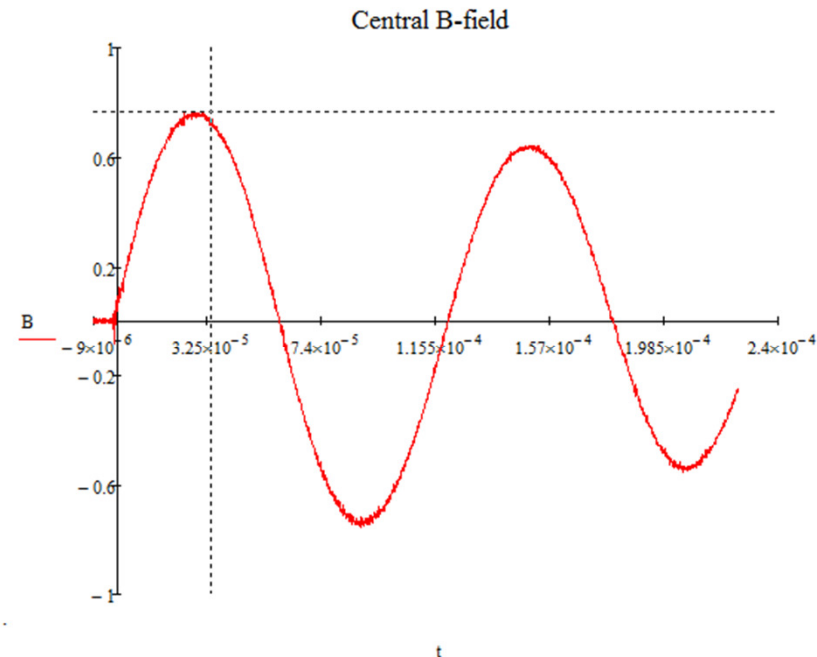
B-field [T]	
FEMM	1.4
Analytical model	1.2
Relative error, %	14.3

The results are close!

10. Field verification using a measuring coil



$$B := \frac{U}{6.4}$$



The B-field measured differs from both analytical and numerical calculation:

B-field [T]	
FEMM	1.4
Analytical model	1.2
Relative error, %	14.3
B measured	0.76...1.9

