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The Laplacian on Cartesian products with mixed boundary conditions

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Abstract. A definition of the Laplacian on Cartesian products with mixed boundary conditions using quadratic forms is proposed. Its consistency with the standard definition for homogeneous and certain mixed boundary conditions is proved and, as a consequence, tensor representations of the corresponding Sobolev spaces of first order are derived. Moreover, a criterion for the domain to belong to the Sobolev space of second order is proved.

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1. Introduction. For $j=1,\ldots,n,\ n\geq 2$, let $\Omega_j\subset\mathbb{R}^{d_j}$ with $d_j\in\mathbb{N}$ be open, and consider $\Omega=\times_{j=1}^n\Omega_j\subset\mathbb{R}^d$ with $d=d_1+\cdots+d_n$. For each j, let \mathcal{H}_j be a closed subspace of $H^1(\Omega_j)$ that is dense in $L^2(\Omega_j)$; possible choices for \mathcal{H}_j are, for instance, $H^1(\Omega_j)$, $H^1_0(\Omega_j)$, $H^1_{\mathrm{per}}(\Omega_j)$ (if applicable), and the subspace of $H^1(\Omega_j)$ corresponding to quasiperiodic boundary conditions (if applicable).

For $f_j: \Omega_j \to \mathbb{C}$, $f_j \in \mathcal{H}_j$, the elementary tensor $f := \bigotimes_{j=1}^n f_j : \Omega \to \mathbb{C}$ denotes the function in $H^1(\Omega)$ given by $f(x_1, \ldots, x_n) = \prod_{j=1}^n f_j(x_j)$ with $x_j \in \Omega_j$. Set

$$\mathcal{K} := \operatorname{span}_{\mathbb{C}} \left\{ \bigotimes_{j=1}^{n} f_j \colon f_j \in \mathcal{H}_j \right\} \subset H^1(\Omega)$$

and

$$\mathcal{H} := \bigotimes_{j=1}^{n} \mathcal{H}_{j} := \overline{\mathcal{K}}^{H^{1}(\Omega)} \subset H^{1}(\Omega),$$

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where the bar denotes the closure in $H^1(\Omega)$. For a more detailed discussion of tensor products of Hilbert spaces, the reader is referred to [7, Sect. 7.5], [9, Sect. 3.4], and [8, Sect. II.4].

Consider on $L^2(\Omega)$ the form \mathfrak{t} with $Dom[\mathfrak{t}] = \mathcal{H}$ and

$$\mathfrak{t}[f,g] = \sum_{k=1}^{d} \langle \partial_k f, \partial_k g \rangle_{L^2(\Omega)}, \quad f,g \in \mathcal{H}.$$

This form is closed, nonnegative, and densely defined. Thus, there exists a unique (nonnegative) self-adjoint operator T on $L^2(\Omega)$ with $Dom(T) \subset \mathcal{H}$ and

$$\mathfrak{t}[f,g] = \langle Tf,g \rangle_{L^2(\Omega)}$$

for all $f \in \text{Dom}(T)$ and $g \in \mathcal{H}$.

The purpose of the present short note is to propose the above definition of T as the definition of the negative Laplacian $-\Delta$ on $L^2(\Omega)$ with mixed boundary conditions corresponding to the ones encoded in the \mathcal{H}_j . To this end, we show that this way to define T is consistent with the definition of the Laplacian on Ω with homogeneous Dirichlet, Neumann, and (quasi-)periodic and also certain mixed boundary conditions. As a byproduct, this also establishes the tensor representations $H^1_0(\Omega) = \bigotimes_{j=1}^n H^1_0(\Omega_j)$, $H^1(\Omega) = \bigotimes_{j=1}^n H^1(\Omega_j)$, and $H^1_{\text{per}}(\Omega) = \bigotimes_{j=1}^n H^1_{\text{per}}(\Omega_j)$. The latter seem to be folklore, but the present author is not aware of any suitable reference in the literature for the general setting above; the case $\Omega = \mathbb{R}^d$ is clear, whereas the recent work [2] establishes the first two representations (for n=2) only for certain bounded domains such as bounded Lipschitz domains.

With the above notations, it is easy to see that

$$\mathcal{H} = \left(\bigotimes_{j=1}^{n-1} \mathcal{H}_j\right) \otimes \mathcal{H}_n,\tag{1.1}$$

so that we may restrict our considerations to the case n=2 by induction. The rest of this note is now organized as follows:

In Sect. 2, we give an alternative characterization for T in the case n=2 using tensor products of operators which makes the operator more accessible for the remaining considerations. Sect. 3 then proves the consistency of the above definition with the case of standard homogeneous boundary conditions, and tensor representations for the corresponding Sobolev spaces of first order are derived. Also compatible mixed boundary conditions are considered here. Finally, Sect. 4 is devoted to a criterion under which the domain of the operator T belongs to $H^2(\Omega)$. As an example, we discuss the particular case where each Ω_j is one dimensional.

2. An alternative characterization using tensor products of operators. Throughout this section, we consider the case n=2, so that $\Omega=\Omega_1\times\Omega_2$ and $\mathcal{H}=\mathcal{H}_1\otimes\mathcal{H}_2$.

For $j \in \{1, 2\}$, let Δ_j be the Laplacian on Ω_j with form domain \mathcal{H}_j , that is, $-\Delta_j$ is the unique nonnegative self-adjoint operator on $L^2(\Omega_j)$ satisfying $\text{Dom}(\Delta_j) \subset \mathcal{H}_j$ and

$$\sum_{k=1}^{d_j} \langle \partial_k u, \partial_k v \rangle_{L^2(\Omega_j)} = \langle (-\Delta_j) u, v \rangle_{L^2(\Omega_j)}$$

for all $u \in \text{Dom}(\Delta_j)$ and $v \in \mathcal{H}_j$. Consider the operator S on $L^2(\Omega)$ defined by

$$S := (-\Delta_1) \otimes I_2 + I_1 \otimes (-\Delta_2) =: S_1 + S_2,$$

where I_j denotes the identity operator on $L^2(\Omega_j)$. According to [7, Theorem 7.23 and Exercise 7.17 a], this operator is self-adjoint and nonnegative with operator core $\mathcal{D} := \operatorname{span}_{\mathbb{C}}\{f_1 \otimes f_2 \colon f_j \in \operatorname{Dom}(\Delta_j)\}$; cf. also [9, Theorem 8.33], [8, Theorem VIII.33] and [5, Proposition A.2].

The above gives an alternative definition of the Laplacian on Ω with mixed boundary conditions corresponding to the \mathcal{H}_j :

Lemma 2.1. For n = 2, we have T = S.

Proof. Let $f = f_1 \otimes f_2 \in \mathcal{D}$ and $g = g_1 \otimes g_2 \in \mathcal{K}$. By Fubini's theorem, we then have

$$\mathfrak{t}[f,g] = \sum_{j=1}^{2} \sum_{k=1}^{d_j} \langle \partial_k f_j, \partial_k g_j \rangle_{L^2(\Omega_j)} \prod_{l \neq j} \langle f_l, g_l \rangle_{L^2(\Omega_l)}$$

$$= \langle (-\Delta_1) f_1, g_1 \rangle_{L^2(\Omega_1)} \langle f_2, g_2 \rangle_{L^2(\Omega_2)} + \langle f_1, g_1 \rangle_{L^2(\Omega_1)} \langle (-\Delta_2) f_2, g_2 \rangle_{L^2(\Omega_2)}$$

$$= \langle S_1 f, g \rangle_{L^2(\Omega)} + \langle S_2 f, g \rangle_{L^2(\Omega)} = \langle S f, g \rangle_{L^2(\Omega)}.$$

By sesquilinearity and denseness of K in \mathcal{H} with respect to the H^1 -norm, this yields

$$\mathfrak{t}[f,g] = \langle Sf,g \rangle_{L^2(\Omega)}$$

for all $f \in \mathcal{D} \subset \mathcal{H}$ and $g \in \mathcal{H}$. Hence, $\mathcal{D} \subset \mathrm{Dom}(T)$ and Tf = Sf for all $f \in \mathcal{D}$, that is, $S|_{\mathcal{D}} \subset T$. Since \mathcal{D} is a core for S and both S and T are self-adjoint, we conclude that $S = \overline{S|_{\mathcal{D}}} = T$.

The above can be considered as a technical preliminary as it makes the operator T more accessible for some of our purposes.

- 3. Consistency with homogeneous and mixed boundary conditions of product type. In this section, we show that the definition of T (with general $n \geq 2$) agrees with the usual definition of the Laplacian in the case where homogeneous boundary conditions (Dirichlet, Neumann, (quasi-)periodic) are imposed on Ω , that is, where all \mathcal{H}_j are chosen of the same type. We also discuss the case of certain mixed boundary conditions compatible with the product structure.
- **3.1. Dirichlet boundary conditions.** We denote the Dirichlet Laplacian on Ω by Δ^D_Ω .

Proposition 3.1. If $\mathcal{H}_j = H_0^1(\Omega_j)$ for each j, then $T = -\Delta_{\Omega}^D$. In particular, we have $H_0^1(\Omega) = \bigotimes_{j=1}^n H_0^1(\Omega_j)$.

Proof. In view of (1.1), it suffices to consider the case n=2 by induction.

Let $f = f_1 \otimes f_2$ with $f_j \in \text{Dom}(\Delta_j)$, and let $g \in C_c^{\infty}(\Omega)$. We then have $h := g(\cdot, x_2) \in C_c^{\infty}(\Omega_1)$ for every $x_2 \in \Omega_2$, and, with S_1 as in Sect. 2 and using Fubini's theorem, we obtain

$$\langle S_1 f, g \rangle_{L^2(\Omega)} = \int_{\Omega_2} f_2(x_2) \langle (-\Delta_1) f_1, h \rangle_{L^2(\Omega_1)} \, \mathrm{d}x_2$$

$$= \sum_{k=1}^{d_1} \int_{\Omega_2} f_2(x_2) \langle \partial_k f_1, \partial_k h \rangle_{L^2(\Omega_1)} \, \mathrm{d}x_2$$

$$= \sum_{k=1}^{d_1} \langle \partial_k f, \partial_k g \rangle_{L^2(\Omega)}.$$

In the same way we see that

$$\langle S_2 f, g \rangle_{L^2(\Omega)} = \sum_{k=d_1+1}^{d_1+d_2} \langle \partial_k f, \partial_k g \rangle_{L^2(\Omega)},$$

and summing up gives

$$\langle Sf, g \rangle_{L^2(\Omega)} = \sum_{k=1}^d \langle \partial_k f, \partial_k g \rangle_{L^2(\Omega)}.$$

By sesquilinearity and denseness of $C_c^{\infty}(\Omega)$ in $H_0^1(\Omega)$, the latter extends to all $f \in \mathcal{D} \subset H_0^1(\Omega_1) \otimes H_0^1(\Omega_2) \subset H_0^1(\Omega)$ and $g \in H_0^1(\Omega)$, which implies that $S|_{\mathcal{D}} \subset -\Delta_{\Omega}^D$. Since \mathcal{D} is a core for S and S and $-\Delta_{\Omega}^D$ are self-adjoint, we conclude that $S = \overline{S|_{\mathcal{D}}} = -\Delta_{\Omega}^D$. In view of Lemma 2.1, this completes the proof of the identity $T = -\Delta_{\Omega}^D$.

Finally, the claimed tensor representation of $H_0^1(\Omega)$ follows since with $T = -\Delta_{\Omega}^D$, also the form domains of T and $-\Delta_{\Omega}^D$ agree.

3.2. Neumann boundary conditions. We denote the Neumann Laplacian on Ω by Δ_{Ω}^{N} .

Proposition 3.2. If $\mathcal{H}_j = H^1(\Omega_j)$ for each j, then $T = -\Delta_{\Omega}^N$. In particular, we have $H^1(\Omega) = \bigotimes_{j=1}^n H^1(\Omega_j)$.

Proof. It again suffices to consider the case n=2.

For $g \in H^1(\Omega)$, we have $g(\cdot, x_2) \in H^1(\Omega_1)$ for almost every $x_2 \in \Omega_2$ and $g(x_1, \cdot) \in H^1(\Omega_2)$ for almost every $x_1 \in \Omega_1$. As in the proof of Proposition 3.1, we therefore obtain

$$\langle Sf, g \rangle_{L^2(\Omega)} = \sum_{k=1}^d \langle \partial_k f, \partial_k g \rangle_{L^2(\Omega)}$$

for all $f \in \mathcal{D} \subset H^1(\Omega_1) \otimes H^1(\Omega_2) \subset H^1(\Omega)$ and all $g \in H^1(\Omega)$; here, we do not need to approximate g but may directly work with $g \in H^1(\Omega)$. The rest of the proof then works exactly as in the above proof.

3.3. Periodic boundary conditions. If Ω is a hyperrectangle, we denote the corresponding Laplacian with periodic boundary conditions by $\Delta_{\Omega}^{\text{per}}$.

Proposition 3.3. For each $j=1,\ldots,n$, let Ω_j be a hyperrectangle in \mathbb{R}^{d_j} and $\mathcal{H}_j=H^1_{\mathrm{per}}(\Omega_j):=\overline{C^\infty_{\mathrm{per}}(\Omega_j)}^{H^1(\Omega_j)}$. Then $T=-\Delta^{\mathrm{per}}_\Omega$. In particular, we have $H^1_{\mathrm{per}}(\Omega)=\bigotimes_{j=1}^n H^1_{\mathrm{per}}(\Omega_j)$.

Proof. For n=2 and $g \in C^{\infty}_{per}(\Omega)$, we have $g(\cdot, x_2) \in C^{\infty}_{per}(\Omega_1)$ for all $x_2 \in \Omega_2$ and $g(x_1, \cdot) \in C^{\infty}_{per}(\Omega_2)$ for all $x_1 \in \Omega_1$. We may now proceed in the same way as in the proof of Proposition 3.1.

Quasiperiodic boundary conditions work analogously.

3.4. Mixed boundary conditions of product type. For each j, let Γ_j be a closed subset of $\partial\Omega_j$, and define the closed subset $\Gamma\subset\partial\Omega$ by

$$\Gamma := (\Gamma_1 \times \overline{\Omega}_2 \times \cdots \times \overline{\Omega}_n) \cup \cdots \cup (\overline{\Omega}_1 \times \cdots \times \overline{\Omega}_{n-1} \times \Gamma_n).$$

Let

$$V := \overline{\{u|_{\Omega} \colon u \in C_c^{\infty}(\mathbb{R}^d \setminus \Gamma)\}}^{H^1(\Omega)},$$

and denote by Δ_{Ω}^{V} the Laplacian on Ω with form domain V. Following [6, Sect. 4.1], this construction can be interpreted as to encode Dirichlet boundary conditions on Γ and Neumann boundary conditions on the rest of the boundary.

Proposition 3.4. If $\mathcal{H}_j = \overline{\{u|_{\Omega_j} : u \in C_c^{\infty}(\mathbb{R}^{d_j} \setminus \Gamma_j)\}}^{H^1(\Omega_j)}$ for each j, then $T = -\Delta_{\Omega}^V$ and $V = \bigotimes_{j=1}^n \mathcal{H}_j$.

Proof. It again suffices to consider the case n=2.

For $u \in C_c^{\infty}(\mathbb{R}^d \setminus \Gamma)$, it is not hard to see that $u(\cdot, x_2) \in C_c^{\infty}(\mathbb{R}^{d_1} \setminus \Gamma_1)$ for all $x_2 \in \Omega_2$ and $u(x_1, \cdot) \in C_c^{\infty}(\mathbb{R}^{d_2} \setminus \Gamma_2)$ for all $x_1 \in \Omega_1$. Since also $(\mathbb{R}^{d_1} \setminus \Gamma_1) \times (\mathbb{R}^{d_2} \setminus \Gamma_2) \subset \mathbb{R}^d \setminus \Gamma$ and, therefore, $\mathcal{H}_1 \otimes \mathcal{H}_2 \subset V$, we may now proceed just as before to conclude the claim.

For n=2 and the choices $\Gamma_1=\partial\Omega_1$ and $\Gamma_2=\emptyset$, we have $\mathcal{H}_1=H^1_0(\Omega_1)$ and, provided that Ω_2 has the so-called *segment property*, $\mathcal{H}_2=H^1(\Omega_2)$, see, e.g., [1, Theorem 3.22]; cf. also [3, Corollary 9.8]. For this situation, Proposition 3.4 shows that the proposed definition of the operator T is consistent with the established one with this kind of mixed boundary conditions. However, if Ω_2 does not have the segment property and $\overline{\{u|_{\Omega_2}: u \in C_c^\infty(\mathbb{R}^{d_2})\}}^{H^1(\Omega_2)} \subsetneq H^1(\Omega_2)$, for the choices $\mathcal{H}_1=H^1_0(\Omega_1)$ and $\mathcal{H}_2=H^1(\Omega_2)$, the present author is not aware of any established way of defining the operator to compare with.

4. Domain in H^2 and an example. In this last section, we prove a criterion in terms of the Δ_j under which the domain of T belongs to $H^2(\Omega)$.

Proposition 4.1. Suppose that for $j=1,\ldots,n,$ we have $\mathrm{Dom}(\Delta_j)\subset H^2(\Omega)$ with

$$\sum_{|\alpha|=2} \frac{1}{\alpha!} \langle \partial^{\alpha} u, \partial^{\alpha} v \rangle_{L^{2}(\Omega_{j})} = \frac{1}{2} \langle \Delta_{j} u, \Delta_{j} v \rangle_{L^{2}(\Omega_{j})}$$
(4.1)

for all $u, v \in \text{Dom}(\Delta_j)$. Then also $\text{Dom}(T) \subset H^2(\Omega)$ with

$$\sum_{|\alpha|=2} \frac{1}{\alpha!} \langle \partial^{\alpha} f, \partial^{\alpha} g \rangle_{L^{2}(\Omega)} = \frac{1}{2} \langle Tf, Tg \rangle_{L^{2}(\Omega_{j})}$$

for all $f, g \in Dom(T)$.

Proof. By induction, it again suffices to consider the case n=2.

Let $f = f_1 \otimes f_2, g = g_1 \otimes g_2 \in \mathcal{D} \subset H^2(\Omega)$. Then, with the notations from Sect. 2,

$$\langle \Delta_1 f_1, \Delta_1 g_1 \rangle_{L^2(\Omega_1)} \langle f_2, g_2 \rangle_{L^2(\Omega_2)} = \langle S_1 f, S_1 g \rangle_{L^2(\Omega)},$$

as well as

$$\langle f_1, g_1 \rangle_{L^2(\Omega_1)} \langle \Delta_2 f_2, \Delta_2 g_2 \rangle_{L^2(\Omega_2)} = \langle S_2 f, S_2 g \rangle_{L^2(\Omega)}$$

and

$$\langle f_1, \Delta_1 g_1 \rangle_{L^2(\Omega_1)} \langle f_2, \Delta_2 g_2 \rangle_{L^2(\Omega_2)} = \langle S_1 f, S_2 g \rangle_{L^2(\Omega)} = \langle S_2 f, S_1 g \rangle_{L^2(\Omega)}.$$

Writing $\alpha = (\beta, \gamma)$ for multiindices in $\mathbb{N}_0^d = \mathbb{N}_0^{d_1} \times \mathbb{N}_0^{d_2}$, we therefore obtain from the hypotheses that

$$\sum_{|\alpha|=2} \frac{1}{\alpha!} \langle \partial^{\alpha} f, \partial^{\alpha} g \rangle_{L^{2}(\Omega)}$$

$$= \sum_{m=0}^{2} \sum_{|\beta|=m} \frac{1}{\beta!} \langle \partial^{\beta} f_{1}, \partial^{\beta} g_{1} \rangle_{L^{2}(\Omega_{1})} \sum_{|\gamma|=2-m} \frac{1}{\gamma!} \langle \partial^{\gamma} f_{2}, \partial^{\gamma} g_{2} \rangle_{L^{2}(\Omega_{2})}$$

$$= \frac{1}{2} \langle S_{1} f, S_{1} g \rangle_{L^{2}(\Omega)} + \langle S_{1} f, S_{2} g \rangle_{L^{2}(\Omega)} + \frac{1}{2} \langle S_{2} f, S_{2} g \rangle_{L^{2}(\Omega)}$$

$$= \frac{1}{2} \langle S f, S g \rangle_{L^{2}(\Omega)} = \frac{1}{2} \langle T f, T g \rangle_{L^{2}(\Omega)}.$$

The latter extends by sesquilinearity to all $f, g \in \mathcal{D}$. In turn, since \mathcal{D} is an operator core for S = T, it extends to all $f, g \in \text{Dom}(T)$ by approximation. In particular, we have $\text{Dom}(T) \subset H^2(\Omega)$. This completes the proof.

Condition (4.1) usually entails that Green's formula can be justified for the partial derivatives of first order with an overall vanishing sum of boundary terms. Examples of domains Ω_j where this can be done for $\mathcal{H}_j = H_0^1(\Omega_j)$ and $\mathcal{H}_j = H^1(\Omega_j)$ are discussed in a slightly different context in a recent joint work with M. Egidi, see [4, Proposition 2.10 and Remark 2.4]. In the present note, the above criterion is demonstrated just for the case where each Ω_j is one dimensional:

Example 4.2. For each j, suppose that $\Omega_j \subset \mathbb{R}$ and that \mathcal{H}_j is such that $\mathrm{Dom}(\Delta_j) \subset H^2(\Omega_j)$ with $\Delta_j u = u''$ for all $u \in \mathrm{Dom}(\Delta_j)$. Then (4.1) is clearly satisfied, and from Proposition 4.1 it follows that $\mathrm{Dom}(T) \subset H^2(\Omega)$ with

$$\sum_{|\alpha|=2} \frac{1}{\alpha!} \langle \partial^{\alpha} f, \partial^{\alpha} g \rangle_{L^{2}(\Omega)} = \frac{1}{2} \langle Tf, Tg \rangle_{L^{2}(\Omega_{j})}$$

for all $f,g \in \text{Dom}(T)$. Note that $\text{Dom}(\Delta_j) \subset H^2(\Omega_j)$ with $\Delta_j u = u''$ for all $u \in \text{Dom}(\Delta_j)$ certainly holds if \mathcal{H}_j contains $C_c^{\infty}(\Omega_j)$, which covers the standard choices for \mathcal{H}_j corresponding to Dirichlet, Neumann, or—if Ω_j is bounded—mixed or (quasi-)periodic boundary conditions.

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